

# Dyadics

A dyadic is the result of a product between two vectors:

$$\underline{D}_s = \underline{u}_s \underline{v}_s$$

This operation provides a description for important properties. For example, we will see that the inertia matrix of a rigid body can be represented by a dyadic. Thus, the dynamics equations can be written in a vector form.

In the following, we will see some properties of the dyadics.

## Properties

1) This product is not commutative.

$$\underline{u}_s \underline{v}_s \neq \underline{v}_s \underline{u}_s$$

2) The dot product between a vector and a dyadic is a vector:

$$\underline{w}_s = \underline{h}_s \cdot \underline{D}_s$$

$\hookrightarrow \underline{h}_s \cdot \underline{D}_s = \boxed{\underline{h}_s \cdot \underline{u}_s} \underline{v}_s = \overset{\text{scalar} = \lambda}{\lambda} \underline{v}_s = \underline{w}_s$

This operation is also not commutative:

$$\underline{h}_s \cdot \underline{D}_s \neq \underline{D}_s \cdot \underline{h}_s$$

$\hookrightarrow \underline{h}_s \cdot \underline{D}_s = \boxed{\underline{h}_s \cdot \underline{u}_s} \underline{v}_s = \overset{\text{scalar} = \lambda}{\lambda} \underline{v}_s$

$\hookrightarrow \underline{D}_s \cdot \underline{h}_s = \underline{u}_s \underline{v}_s \cdot \underline{h}_s = \overset{\text{scalar} = \alpha}{\alpha} \underline{u}_s$

3) We have the following definitions:

$$\underline{J}_s = (\underline{w}_s \times \underline{u}_s) \underline{v}_s = \underline{w}_s \times \underline{u}_s \underline{v}_s \triangleq \underline{w}_s \times \underline{D}_s$$

$$\underline{K}_s = \underline{u}_s (\underline{v}_s \times \underline{w}_s) = \underline{u}_s \underline{v}_s \times \underline{w}_s \triangleq \underline{D}_s \times \underline{w}_s$$

In general,  $\underline{J}_s \neq \underline{K}_s$ .

## Matrix representation of dyadics

Let  $\underline{u}_s$  and  $\underline{v}_s$  be two vectors in  $\mathbb{R}^3$  with representations  $u_a$  and  $v_a$  in the basis  $\underline{e}_a$ . Thus:

$$\underline{D}_s = \underline{u}_s \underline{v}_s = \underline{e}_a^T u_a v_a^T \underline{e}_a = \underline{e}_a^T \underline{D}_a \underline{e}_a$$

Thus, the representation of the dyadics is a  $3 \times 3$  matrix such as:

$$\underline{D}_a = \underline{u}_a \underline{v}_a^T$$

The matrix form can be obtained from the vectorial form using:

$$\underline{D}_a = \underline{e}_a \cdot \underline{D}_a \cdot \underline{e}_a^T$$

$$\hookrightarrow \underline{e}_a \cdot \underline{D}_a \cdot \underline{e}_a^T = \underline{e}_a \cdot \underline{u}_a \underline{v}_a^T \cdot \underline{e}_a^T = \underbrace{\underline{e}_a \cdot \underline{e}_a^T}_{=\underline{I}} \underline{u}_a \underline{v}_a^T \underbrace{\underline{e}_a \cdot \underline{e}_a^T}_{=\underline{I}} = \underline{u}_a \underline{v}_a^T = \underline{D}_a$$

The dot product between a dyadic and two vectors results in a scalar:

$$\underline{u}_a \cdot \underline{D}_a \cdot \underline{v}_a = \underline{u}_a^T \underbrace{\underline{e}_a \cdot \underline{e}_a^T}_{=\underline{I}} \underline{D}_a \underbrace{\underline{e}_a \cdot \underline{e}_a^T}_{=\underline{I}} \underline{v}_a = \underline{u}_a^T \underline{D}_a \underline{v}_a$$

$1 \times 3 \quad 3 \times 3 \quad 3 \times 1 = 1 \times 1 \text{ (scalar)}$

A null dyadics is a dyadic in which its representation is a null matrix in  $\mathbb{R}^{3 \times 3}$ :

$$\underline{D}_a = \underline{0}_a \Rightarrow \underline{D}_a = \underline{0}_{3 \times 3}$$

A unitary dyadics is a dyadic in which its representation is an identity matrix in  $\mathbb{R}^{3 \times 3}$ :

$$\underline{D}_a = \underline{1}_a \Rightarrow \underline{D}_a = \underline{I}_{3 \times 3}$$

In this case:

$$\underline{u}_a = \underline{u}_a \underline{1}_a = \underline{1}_a \underline{u}_a$$

$$\hookrightarrow \underline{e}_a^T \underline{I} \underline{e}_a \underline{u}_a = \underline{e}_a^T \underline{u}_a = \underline{u}_a$$

## Representation of dyadics in different basis

Given a dyadic  $\underline{D}_a$  with representation  $\underline{D}_a$  in the basis  $\underline{e}_a$ , we can find its representation in the basis  $\underline{e}_b$  using:

$$\underline{D}_a = \underline{e}_a \cdot \underline{D}_a \cdot \underline{e}_a^T \Rightarrow \underline{e}_a^T \underline{D}_a \underline{e}_a = \underline{D}_a$$

$$\underline{D}_b = \underline{e}_b \cdot \underline{D}_a \cdot \underline{e}_b^T$$

$$\underline{D}_b = \underbrace{\underline{e}_b \cdot \underline{e}_a^T}_{\underline{C}_{ba}} \underline{D}_a \underbrace{\underline{e}_a \cdot \underline{e}_b^T}_{\underline{C}_{ab} = \underline{C}_{ba}^T} = \underline{C}_{ba} \underline{D}_a \underline{C}_{ab} = \underline{C}_{ba} \underline{D}_a \underline{C}_{ba}^T$$

$$\underline{D}_b = \underline{C}_{ba} \underline{D}_a \underline{C}_{ab} = \underline{C}_{ba} \underline{D}_a \underline{C}_{ba}^T$$

If the vectors  $\underline{u}_a$  and  $\underline{v}_a$  are known in different basis, then:

$$\underline{D}_a = \underline{u}_a \underline{v}_a^T = \underline{e}_a^T \underline{u}_a \underline{v}_a^T \underline{e}_b = \underline{e}_a^T \underline{D}_{ab} \underline{e}_b$$

Analogously:

$$\underline{D}_{ab} = \underline{u}_a \underline{v}_b^T = \underline{e}_a \cdot \underline{u}_a \underline{v}_b^T \cdot \underline{e}_b^T = \underline{e}_a \cdot \underline{D}_a \cdot \underline{e}_b^T$$

If it is desired to represent the dyadic entirely in a single basis, then:

$$D_a = \underline{e}_a \cdot D_s \cdot \underline{e}_a^T = \underbrace{\underline{e}_a \cdot \underline{e}_a^T}_I D_{ab} \underbrace{\underline{e}_b \cdot \underline{e}_a^T}_{C_{ba} = C_{ab}^T} = D_{ab} C_{ab}^T$$

$$D_b = \underline{e}_b \cdot D_s \cdot \underline{e}_b^T = \underbrace{\underline{e}_b \cdot \underline{e}_a^T}_{C_{ba} = C_{ab}^T} D_{ab} \underline{e}_a \cdot \underline{e}_a^T = C_{ab}^T D_{ab}$$